# HERMES spectroscopy of normal A and Am stars 

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#### Abstract

The nominal Kepler mission provided very high precision photometric data. Using these data, interesting phenomena such as spots, and 'hump and spike' features were observed in the light curves of some normal A and metallic lined A stars (Am stars). However, the connection between such phenomena and the chemical peculiarity of the Am stars is still unclear. In order to make progress on these issues, it's important to collect high-resolution spectroscopic data to determine their fundamental parameters and individual chemical abundances. In this paper, we present a spectroscopic study of a sample of 'hump and spike' stars in the nominal Kepler field. We used data collected with the High Efficiency and Resolution Mercator Échelle Spectrograph (HERMES). We determined the spectral types of these stars and obtained the atmospheric stellar parameters such as effective temperatures, surface gravities, projected rotational, microturbulent, and radial velocities. We also performed a detailed individual chemical abundance analysis for each target. We confirmed KIC 3459226 and KIC 6266219 as Am stars, KIC 9349245 as a marginal Am star, while KIC 4567097, KIC 4818496, KIC 5524045, KIC 5650229, KIC 7667560, and KIC 9272082 are non-Am stars. To estimate their evolutionary phases, all the stars were placed in the Hertzsprung-Russell diagram. Based on their spectral classification and chemical abundance pattern, we reclassified KIC 6266219 (previously treated as chemically normal) as an Am star (kA3hA7mF1) and KIC 9272082 (previously treated as Am) as non-Am.


Key words: techniques: spectroscopic-stars: abundances-stars: chemically peculiar-stars: general-stars: rotationstarspots.

## 1 INTRODUCTION

Over the previous few decades, many space missions were lauched, such as Microvariability and Oscillations of Stars (MOST; Rucinski et al. 2003), Convection, Rotation and planetary Transits (CoRoT; Baglin et al. 2006), Kepler (Borucki et al. 2010), and Transiting Exoplanet Survey Satellite (TESS; Ricker et al. 2014), which are mainly searching for exoplanets and enabling asteroseismology of pulsating stars. More projects like PLAnetary Transits and Oscillations of stars (PLATO; Rauer et al. 2014) and Wide-Field Infrared Survey Telescope (WFIRST; Gehrels, Spergel \& WFIRST SDT Project 2015) will be launched in the near future. To understand the characteristics of exoplanets, the stellar parameters of the host stars are relevant (Howard et al. 2013). For pulsating stars, especially those with solar-like oscillations, asteroseismology has played a big role in determining the stellar parameters using high-precision photometric data from space telescopes. Even when high-precision photometric data are available, ground-based spectroscopy is still essential in determining complementary stellar parameters of most stars and it is still widely used (e.g. Santos et al. 2005; Sousa et al. 2015; Gandolfi et al. 2018; Kassounian et al. 2019; Eisner et al. 2020). For the intermediate-mass main-sequence (MS) stars (A- and F-type), where chemically peculiar (CP) stars are common (about 10 per cent),

[^0]spectroscopy plays a big role in determining the individual chemical abundances in addition to the atmospheric stellar parameters (e.g. Fossati et al. 2008; Gebran \& Monier 2008; Gebran, Monier \& Richard 2008; Gebran et al. 2010).

The CP stars exhibit atmospheric abundances that are significantly different from the solar values. Based on the magnetic field strength and strength of absorption lines in their optical spectra, Preston (1974) categorized the CP stars into four major groups: CP1 stars (the metallic-line or Am/Fm stars), CP2 stars (the magnetic Ap stars), CP3 stars (the Mercury-Manganese or HgMn stars), and CP4 stars (the He-weak stars). The CP1 (Am) are the most numerous CP stars followed by CP2 (Ap) stars in the intermediate-mass MS part of the Hertzsprung-Russell (HR) diagram. The Am stars are characterized by underabundance of certain light elements (such as Ca and Sc ) and excess of metals, in addition to the very weak or non-detectable magnetic fields in general (Conti 1970; Preston 1974; Romanyuk 2007). This implies that for Am stars, three spectral classifications are always given: (i) based on the Balmer lines, which give good estimates of effective temperature ( $T_{\text {eff }}$ ), (ii) based on the Ca II Kline ( $\lambda \lambda 393.4 \mathrm{~nm}$ ), which gives an earlier spectral type since it is weaker than for normal stars, and (iii) based on the metallic lines, which give a later spectral type because they are enhanced relative to normal stars. Am stars whose spectral types are derived from the CaII K-line and metallic lines differ by five or more subtypes are called classical Am stars. For marginal Am stars, this difference is less than five subtypes. The CP2 stars are known for large excesses


Figure 1. The 'hump and spike' features in an amplitude spectrum for KIC 6192566 as reported by Trust et al. (2020).
(up to several orders of magnitude) of heavy elements such as $\mathrm{Si}, \mathrm{Hg}$ or the rare-Earth elements $\mathrm{Sr}, \mathrm{Cr}, \mathrm{Eu}, \mathrm{Nd}$, Pr etc (Conti 1970; Preston 1974; Kurtz \& Martinez 2000; Romanyuk 2007). The CP2 stars are also characterized by the presence of well organized magnetic fields with strengths of up to several tens of kG (Aurière et al. 2007), which could be of fossil origin (Braithwaite \& Spruit 2004). The CP3 stars possess enhanced $\mathrm{Hg}_{\text {II }}$ ( $\lambda \lambda 398.4 \mathrm{~nm}$ ) and/or Mn II lines and weak lines of light elements (such as $\mathrm{He}, \mathrm{Al}$, and N ) in their spectra. The majority of CP3 stars are very slow rotators with $v \sin i$ $\leq 20 \mathrm{~km} \mathrm{~s}^{-1}$ and are also characterized by weak or non-detectable magnetic fields (Preston 1974; Smith 1997; Bohlender, Dworetsky \& Jomaron 1998). The CP4 stars are characterized by anomalously weak He I lines in their spectra. They are also slow rotators with magnetic field strengths of upto the orders of 1 kG (Jones \& Wolff 1974; Preston 1974).

The chemical peculiarities in CP stars are thought to mainly originate from the interplay between radiative levitation and gravitational settling, a process known as atomic diffusion (Michaud 1970; Watson 1970; Khokhlova 1981; Hui-Bon-Hoa 2000; Richer, Michaud \& Turcotte 2000; Turcotte 2003; Théado et al. 2011). Based on this theory, in the absence of mixing, light elements drift under the influence of gravity and are seen as underabundant while heavy elements are radiatively driven outward (reflected as overabundant). This theory requires calm and stable atmospheres that are facilitated by the slow rotation of CP stars (Fossati et al. 2008; Takeda et al. 2008; Abt 2009; Stateva et al. 2009) and by strong magnetic fields in the case of Ap stars. In addition to turbulent mixing, another actor in the process could be weak mass-loss (Richer et al. 2000; Richard, Michaud \& Richer 2001; Michaud et al. 2004; Talon, Richard \& Michaud 2006). It has also been reported to reduce the main surface chemical abundance anomalies, in models including atomic diffusion, to the observed levels in CP stars (Michaud et al. 1983; Michaud \& Charland 1986; Alecian 1996; Vick et al. 2010, 2011). The strong magnetic fields are thought to stabilize the convective material, while slow rotation minimizes meridional circulation and in-turn minimizes mixing that would counteract atomic diffusion. Rotation plays a big role in constraining the physics of CP stars such as pulsations (Dziembowski, Krolikowska \& Kosovichev 1988; Soufi, Goupil \& Dziembowski 1998), overshooting (Browning, Brun
\& Toomre 2004; Costa et al. 2019), and the observed chemical abundances (Turcotte 2002; Murphy 2014).

Given the high precision and almost continuous photometric data from space missions like the nominal Kepler mission, K2 (Howell et al. 2014), and TESS, the rotation of A-type (CP and normal) stars has been studied assuming spot induced rotational modulation (Kron 1947; Mosser et al. 2009; Balona 2013) and based on the 'hump and spike' features in their frequency spectra (Saio et al. 2018; Trust et al. 2020). Briefly, the 'hump and spike' stars were named by Saio et al. (2018). Their frequency spectra possess a sharp peak ('spike') on the high frequency side of a broad hump of very close frequencies ('hump') as shown in Fig. 1 by Trust et al. (2020). The broad humps are induced by Rossby waves ( r modes) and the spikes are the rotation frequencies induced by one or more spots (Saio et al. 2018). The 'hump and spike' stars are a test bed to study the association of the r modes with rotation and other processes (such as atomic diffusion) that affect their chemical peculiarity. The period of stellar rotation can also be determined using seismic splitting (Claverie et al. 1981; Saio 1981) and chromospheric activity, especially Ca II emission (Noyes et al. 1984). However, we obtain a complete rotational profile when the projected rotational velocity $(v \sin i)$ is known too. The $v \sin i$ can be spectroscopically retrieved from the broadening of the absorption lines (Kaler 1989).

To search and study the pulsational variabilities in Ap and $\mathrm{Am} / \mathrm{Fm}$ stars, a dedicated ground-based project, the Nainital-Cape survey, was initiated between astronomers of India and South Africa. With time, astronomers from other countries joined this programme making it a multinational collaborative project and a number of results have been published (e.g. Ashoka et al. 2000; Martinez et al. 2001; Joshi et al. 2003, 2006, 2009, 2010, 2012, 2016, 2017). The success of the Nainital-Cape survey and the discovery of 'hump and spike' features in the frequency spectra (Balona 2013; Balona et al. 2015; Saio et al. 2018) of a number of normal A and Am stars motivated us to initiate a collaboration with astronomers from India, Uganda, and Belgium, to study the phenomena responsible for the observed chemical peculiarities in Am stars. The first results are presented in Trust et al. (2020). This paper is the second of the series whose objective is to characterize the 'hump and spike' stars. This involves the spectral classification, determination of atmospheric stellar parameters, and individual chemical abundances of our targets.

This paper is organized as follows. The data are discussed in Section 2. Spectral classification is presented in Section 3, while atmospheric parameters and individual chemical abundance analyses are given in Section 4. The comments on individual stars and their location in the HR diagram are discussed in Sections 5 and 6, respectively, while the conclusion is given in Section 7.

## 2 THE DATA

### 2.1 Selection criteria

All the stars in this study are part of the samples analysed by Trust et al. (2020). These are normal A and Am stars with 'hump and spike' features in their frequency spectra which were selected from Balona (2013), Balona et al. (2015), and Gray et al. (2016). Of the 170 'hump and spike' stars studied by Trust et al. (2020), we selected those with the $V \leq 10.5 \mathrm{mag}$ and not reported in the literature as members of binary systems. Stars with $V \leq 10.5 \mathrm{mag}$ are bright enough to obtain high-quality High Efficiency and Resolution Mercator Échelle Spectrograph (HERMES) spectra with signal-to-noise ratio (SNR) of at least 100 per pixel in a maximum exposure time of 1 h . In this paper, we present a spectroscopic analysis for nine such stars.

Table 1. The observation $\log$ for the target stars. The right ascension $\left(\alpha_{\mathrm{J} 2000}\right)$ and declination $\left(\delta_{\mathrm{J} 2000}\right)$ given by Gaia Collaboration (2018) and the $V$-band magnitude estimated by Høg et al. (2000a) are included.

| $\begin{aligned} & \text { KIC } \\ & \text { no. } \end{aligned}$ | Other name | $\begin{gathered} \alpha_{\mathrm{J} 2000} \\ (\mathrm{hh} \mathrm{~mm} \mathrm{ss}) \end{gathered}$ | $\begin{gathered} \delta_{\mathrm{J} 2000} \\ (\mathrm{dd} \mathrm{~mm} \mathrm{ss}) \end{gathered}$ | $\begin{gathered} V \\ (\mathrm{mag}) \end{gathered}$ | BJD (2450000+) <br> (d) | Number of spectra | $\begin{gathered} \text { SNR } \\ \text { (at } 500,650, \\ \text { and } 810 \mathrm{~nm} \text { ) } \end{gathered}$ | Total exposure duration (min) |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3459226 | $\mathrm{BD}+38^{\circ} 3705$ | 194050.1201 | +38 3227.0912 | 10.11 | 8430 | 5 | 27-51 | 75.0 |
| 4567097 | HD 184469 | 193255.9360 | +39 3945.7237 | 7.75 | 8429 | 1 | 90-111 | 10.0 |
| 4818496 | HD 177592 | 190339.2582 | +395439.2370 | 8.07 | 8429 | 1 | 99-125 | 12.5 |
| 5524045 | $\mathrm{BD}+40^{\circ} 3639$ | 191514.7706 | +40 4345.3009 | 9.36 | 8429 | 2 | 42-71 | 40.0 |
| 5650229 | HD 226697 | 195636.7584 | +40 4832.0652 | 9.92 | 8430 | 3 | 37-55 | 42.5 |
| 6266219 | TYC3127-2016-1 | 185658.4055 | +413645.3267 | 9.87 | 8425 | 3 | 25-44 | 60.0 |
| 7667560 | HD 177458 | 190249.8015 | +432121.1749 | 9.63 | 8425 | 3 | 23-51 | 60.0 |
| 9272082 | HD 179458 | 191033.7965 | +454503.9964 | 8.95 | 8429 | 4 | 55-86 | 50.0 |
| 9349245 | HD 185658 | 193805.1898 | +455303.8221 | 8.11 | 8430 | 1 | 70-104 | 10.0 |

Table 2. The spectral classification of the target stars.

| KIC | Sp type | Comment | References |
| :---: | :---: | :---: | :---: |
| 3459226 | F2 Dwarf |  | 1a |
|  | kA2hF0mF3 |  | 1 b |
|  | kA2hF0mF2 |  | 2 |
| 4567097 | A0 Giant |  | 1a |
|  | B9 III | Vgood | 1 b |
|  | B9 |  | 3 |
| 4818496 | A1 Dwarf |  | 1 a |
|  | A1 V | Good | 1 b |
|  | A0 |  | 3 |
| 5524045 | A1 Dwarf |  | 1 a |
|  | A6 mA0 V metal-weak | Good | 1 b |
|  | A3 mA0 IV metal-weak |  | 2 |
|  | A1V |  | 4 |
|  | A0.5Va+ |  | 5 |
| 5650229 | A1 Dwarf |  | 1 a |
|  | A8 mA0 V metal-weak |  | 1 b |
|  | A3 mA1 IV metal-weak |  | 2 |
|  | A0 IV-V |  | 2 |
|  | A2 |  | 3 |
|  | B9 III or IV |  | 4 |
| 6266219 | F1 Dwarf |  | 1 a |
|  | kA3hA7mF1 |  | 1 b |
| 7667560 | A9 Dwarf |  | 1a |
|  | A6 mA0 V metal-weak | Good | 1 b |
|  | A |  | 3 |
| 9272082 | A9 Dwarf |  | 1a |
|  | A4 V | Good | 1 b |
|  | A5 IV-V |  | 2 |
|  | A0 |  | 3 |
|  | A3 V |  | 4 |
|  | A7 |  | 6 |
|  | Am |  | 7 |
|  | A5m |  | 8 |
| 9349245 | F0 Dwarf |  | 1 a |
|  | A8 IV (Sr) | Good | 1 b |
|  | A8 V (Sr) |  | 2 |
|  | A5 |  | 3 |
|  | A7 III |  | 4 |
|  | Am |  | 9 |
|  | F2 III |  | 10 |

Note. References: (1) This study based on: (a) fitting SED with Kesseli et al. (2017) template libraries using the vosA tool, and (b) HERMES spectra using MKCLASS (Gray \& Corbally 2014); (2) Gray et al. (2016); (3) Cannon \& Pickering (1993); (4) Frasca et al. (2016); (5) 'a+' denotes higher luminosity main-sequence star (Niemczura et al. 2015); (6) Floquet (1970); (7) Bertaud (1960); (8) Macrae (1952); (9) Catanzaro et al. (2015); and (10) Niemczura et al. (2017).

### 2.2 Observations and data reduction

The high-resolution spectroscopic observations were made on the nights of 2018 November 2, 6, and 7, using the HERMES (Raskin et al. 2011) spectrograph mounted at the Cassegrain focus of the 1.2-m Mercator telescope located at La Palma, Spain. In a single exposure, this spectrograph records optical spectra in the wavelength ( $\lambda$ ) range $377-900 \mathrm{~nm}$ across the 55 spectral orders. The resolving power of this instrument in the high-resolution mode is 85000 with radial velocity stability of about $50 \mathrm{~m} \mathrm{~s}^{-1}$, and excellent throughput (Raskin et al. 2011). The $\log$ of the observations collected with HERMES is presented in Table 1.

The spectra were reduced using a dedicated HERMES pipeline. The usual reduction procedure for Échelle spectra was applied. The steps include subtraction of the bias and stray light, flat-field correction, order-by-order extraction, wavelength calibration with Thorium-Argon lamps, cosmic rays removal, and merging of the orders (Raskin et al. 2011). The procedure also provides the SNR, which in our case, falls in the range $30-110$ (at $\lambda=500,650$, and 810 nm ) on average and are listed in Table 1. Normalization of all the spectra to the local continuum was performed manually using the continuum task of the IRAF package. ${ }^{1}$ Finally, all the spectra were systematically corrected for barycentric motion.

For some of the stars, we collected more than one spectrum (Table 1). For these stars, all the spectra were median-combined and we performed investigations on the average spectrum. The SNR of the median spectra lies in the range of $110-220$ on average.

## 3 SPECTRAL CLASSIFICATION

Prior to the determination of atmospheric parameters and chemical abundances, we performed a spectral classification analysis of the targets on the MK classification system (Morgan, Keenan \& Kellman 1943; Gray \& Corbally 2009). Spectral classification provides essential details about a star's chemical peculiarity and initial atmospheric parameters (Gray \& Corbally 2009). The spectral type and luminosity class are based on the similarity between the observed spectra and those of well-known standards, taking into account essential hydrogen and metal lines. We classified the stars using two data types: (i) spectral energy distribution (SED) and (ii) HERMES spectra.

Using the VOSA package (Bayo et al. 2008) and least-squares minimization technique, we cross-matched the SED (discussed in

[^1]Table 3. List of basic and photometric parameters. The parallax $(\pi)$ determined by Gaia Collaboration (2018), reddening parameter $[E(B-V)]$ computed using 3D dustmaps by Green et al. (2019), the effective temperature ( $T_{\text {eff }}$ ), surface gravity ( $\log g$ ), and metallicity ( $[\mathrm{M} / \mathrm{H}]$ ) estimated using $u v b y \beta$ indices, $T_{\text {eff }}$ computed from 2MASS, $T_{\text {eff }}$ derived from UBV, and the estimated $T_{\text {eff }}, \log g$, and $[\mathrm{M} / \mathrm{H}]$ from the SED fitting are given.

| KIC | $\pi$ | $E(B-V)$ | $T_{\text {eff }}$ uvby $\beta$ | $\log g$ <br> uvby $\beta$ | $[\mathrm{M} / \mathrm{H}]$ $u v b y \beta$ | $\underset{\text { 2MASS }}{T_{\mathrm{eff}}}$ | $\underset{\text { eff }}{T_{\text {ery }}}$ | $\begin{aligned} & T_{\text {eff }} \\ & \text { SED } \end{aligned}$ | $\underset{\mathrm{SED}}{\log g}$ | $\begin{gathered} {[\mathrm{M} / \mathrm{H}]} \\ \mathrm{SED} \end{gathered}$ |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| No. | (mas) | (mag) | $( \pm 200 \mathrm{~K})$ | $( \pm 0.10$ dex $)$ | $( \pm 0.13 \mathrm{dex})$ | (K) | (K) | $( \pm 250 \mathrm{~K})$ | $( \pm 0.5 \mathrm{dex})$ | (dex) |
| 3459226 | $3.3107 \pm 0.0804$ | 0.0303 | - | - | - | $7390 \pm 280$ | $6990 \pm 250$ | 7500 | 4.0 | $0.0 \pm 0.35$ |
| 4567097 | $2.0990 \pm 0.0480$ | 0.0581 | - | - | - | $9490 \pm 510$ | $10430 \pm 350$ | 9750 | 3.5 | $0.2 \pm 0.25$ |
| 4818496 | $4.0709 \pm 0.0366$ | 0.0593 | - | - | - | $9580 \pm 290$ | $9400 \pm 270$ | 9500 | 4.5 | $0.0 \pm 0.35$ |
| 5524045 | $2.1854 \pm 0.0338$ | 0.0449 | - | - | - | $8900 \pm 360$ | $9170 \pm 260$ | 9250 | 4.5 | $0.0 \pm 0.35$ |
| 5650229 | $1.6616 \pm 0.0286$ | 0.1377 | - | - | - | $9250 \pm 340$ | $9450 \pm 330$ | 9000 | 3.5 | $0.5 \pm 0.30$ |
| 6266219 | $2.5992 \pm 0.0362$ | 0.0310 | - | - | - | $7930 \pm 260$ | $7160 \pm 260$ | 7750 | 3.5 | $0.2 \pm 0.25$ |
| 7667560 | $2.8792 \pm 0.0347$ | 0.0245 | - | - | - | $8310 \pm 290$ | $8340 \pm 220$ | 8000 | 3.5 | $0.2 \pm 0.25$ |
| 9272082 | $2.5738 \pm 0.0386$ | 0.0254 | 8470 | 4.13 | 0.17 | $8010 \pm 170$ | $7920 \pm 130$ | 8000 | 3.5 | $0.5 \pm 0.30$ |
| 9349245 | $6.9544 \pm 0.0343$ | 0.0100 | 8130 | 4.27 | 0.22 | $7780 \pm 170$ | $7690 \pm 120$ | 7500 | 3.5 | $0.5 \pm 0.30$ |

Section 4.2) of the targets with the template libraries built by Kesseli et al. (2017) from SDSS spectra. The resulting spectral types are indicated with ' 1 a' in Table 2. The uncertainty in the spectral type is about three subtypes.

Classification of the HERMES spectra was performed using the code mKCLASS (Gray \& Corbally 2014). Given that the observation time is always limited and not all the standards are always visible at the time of observation of the targets, we used the available standard libraries by Gray et al. (2003). For the observations of these standard libraries, they used the Dark Sky Observatory (DSO) 0.8-m telescope in Northwestern part of North Carolina in combination with the Gray/Miller classification spectrograph having two gratings with 600 and 1200 grooves $\mathrm{mm}^{-1}$, respectively. The standards were observed in the violet-green spectral region at a resolution range of $0.15-$ 0.36 nm per 2 pixels (Gray \& Corbally 2014).

The spectra for the standards and our sample stars are not from the same spectrograph/grating combination. To approximate the specifications of the standards as close as possible, we truncated the wavelength range, rebinned and convolved the HERMES spectra with a Gaussian of appropriate full width at half-maximum (FWHM) determined by the intended final resolution of 0.36 nm per 2 pixels.

The metric-distance technique (LaSala 1994), which relies on a weighted least-squares comparison of the spectra of the program star with those of the MK standard stars, was used (Gray \& Corbally 2014). For each star, the spectral type was determined three times, each time based on different lines, that is, hydrogen lines ( $\mathrm{H} \gamma$ and $\mathrm{H} \sigma$ lines), metal lines, and the Ca II K -line. For a normal star, all the three regions correspond to the same spectral type. However, in the case of the CP stars, such as Am stars, the three regions give different spectral types (Gray \& Corbally 2009). The obtained spectral types and luminosity classes are indicated with ' 1 b ' in Table 2 . The spectral types and luminosity classes of the sample lie in the range B9 to F1 and III to V, respectively.

## 4 ATMOSPHERIC PARAMETERS AND CHEMICAL ABUNDANCES

Photometry and spectroscopy are the traditional methods used to determine atmospheric parameters. In the case of well-studied pulsating stars, asteroseismology has proved to give good results, especially since high-precision photometric data from space missions like CoRoT and Kepler became available. That does not alter the fact that spectroscopy has always remained robust and relevant in providing atmospheric parameters. However, a good initial guess is very essential to attain convergence during the process of atmospheric
parameter determination based on spectral synthesis. Some of the methods from which the initial guesses can be obtained include photometric colour indices (Section 4.1) and SED fitting (Section 4.2).

### 4.1 Atmospheric parameters from photometry

We computed stellar parameters using the available photometric indices from photometric databases, namely, ubvy $\beta$ (Hauck \& Mermilliod 1998), 2MASS (Cutri et al. 2003; Skrutskie et al. 2006), and UBV (Oja 1985). All values were retrieved from the General Catalog of Photometric data ${ }^{2}$ (Mermilliod, Mermilliod \& Hauck 1997). The reddening parameter $E(B-V)$ was computed from 3D models (Green 2018; Green et al. 2019) using the Gaia parallaxes (Gaia Collaboration 2018) and the stellar Galactic coordinates from the SIMBAD database ${ }^{3}$ (Wenger et al. 2000). The resulting values of $E(B-V)$ are reported in column 3 of Table 3.

In the General Catalog of Photometric data, the Strömgren $\beta$ indices are available for KIC 9272082 and KIC 9349245 only. We calculated $T_{\text {eff }}$ and $\log g$ from $u b v y \beta$ indices using a routine originally written by Moon (1985), based on Moon \& Dworetsky (1985) and corrected by Napiwotzki, Schoenberner \& Wenske (1993). We determined the overall metallicity $([\mathrm{M} / \mathrm{H}])$ using the Strömgren $\delta m_{o}$ index for A-type (A0-F0) stars (Smalley 1993). We considered errors (systematic and statistical) of $200 \mathrm{~K}, 0.10$ dex (Napiwotzki et al. 1993), and 0.13 dex (Smalley 1993) in $T_{\text {eff }}, \log g$, and [M/H], respectively. The results based on Strömgren photometry are given in columns 4-6 of Table 3.

Finally, we computed $T_{\text {eff }}$ from the 2MASS and UBV photometric systems. This is based on the strong correlation of $T_{\text {eff }}$ with the $(V-$ $K)_{o}$ (Masana, Jordi \& Ribas 2006) and ( $\left.B-V\right)_{o}$ indices (Sekiguchi \& Fukugita 2000). Since $T_{\text {eff }}$ is weakly correlated with $\log g$ and [M/H], we fixed $\log g=4.0$ dex and $[\mathrm{M} / \mathrm{H}]=0.0$ dex (solar metallicity). For each system, we estimated the uncertainties in $T_{\text {eff }}$ based on the errors of the photometric indices, i.e. $E(B-V)(0.02 \mathrm{mag}),[\mathrm{M} / \mathrm{H}]$ ( 0.10 dex ), and $\log g$ ( 0.10 dex) (Kahraman Aliçavuş et al. 2016). The photometric parameters and their uncertainties are listed in columns 7 and 8 of Table 3. The $T_{\text {eff }}$ (2MASS) and $T_{\text {eff }}$ (UBV) for KIC 4567097 and KIC 6266219 do not agree within $1 \sigma$. In addition, $T_{\text {eff }}(u b v y \beta)$ values for KIC 9272082 and KIC 9349245 seem to be about 500 K higher compared to the values of $T_{\text {eff }}$ (2MASS) and $T_{\text {eff }}$ (UBV). The difference in $T_{\text {eff }}$ could be introduced by the assumption that ( $V$

[^2]

Figure 2. The red filled circles represent the SEDs of the target stars. The blue solid lines represent the spectra fitted to the data using the vOSA tool.
$-K)$ and $(B-V)$ are completely $\log g$ and metallicity independent (Ramírez \& Meléndez 2005). Some of the target stars in the sample are CP and/or variable, which may introduce additional errors in the calibration process (Herdin, Paunzen \& Netopil 2016).

### 4.2 Atmospheric parameters from SED

We can measure stellar parameters from the SED of a star. SEDs are built using spectrophotometry obtained in different wavelengths, ideally from ultraviolet to infrared. For this purpose, we adopted the vosA tool.

In the first step, the tool automatically searches for spectrophotometric observations in different data bases, i.e. the Tycho II (Høg et al. 2000b), 2MASS (Skrutskie et al. 2006), WISE (Wright et al. 2010), and Gaia DR2 (Maíz Apellániz \& Weiler 2018) catalogues, complemented with Johnson (Johnson \& Morgan 1953; Oja 1985), Strömgren (Hauck \& Mermilliod 1998), Galaxy Evolution Explorer (GALEX; Bianchi \& GALEX Team 2000), Sloan (Brown et al. 2011), and Pan-STARRS1 (Chambers et al. 2016) photometry, when available. We estimated the extinction $\left(A_{\mathrm{v}}\right)$ for each target using $A_{\mathrm{V}}=R_{\mathrm{V}} \times E(B-V)=3.1 \times E(B-V)$ (Brown et al. 2011). We determined $E(B-V)$ as discussed in Section 4.1. Also the stellar distance is required to construct SEDs. We used distances calculated from the Gaia parallaxes (Gaia Collaboration 2018) that are listed in column 2 of Table 3. The SEDs for all the stars under study are shown in Fig. 2 as filled circles. Using ATLAS9 Kurucz ODFNEW/NOVER models (Castelli, Gratton \& Kurucz 1997) in the vOSA tool, we performed a least-squares fit to these SEDs to obtain $T_{\text {eff }}, \log g$, and $[\mathrm{M} / \mathrm{H}]$. The results are represented by the blue lines in Fig. 2. The uncertainties calculated by the vosa tool on these parameters are underestimated by a factor of about 2 because they only account for the internal errors on the SED-fitting procedure. This is attributed to the intrinsic limitations of the method (Bayo et al. 2008; McDonald, Zijlstra \& Boyer 2012; Catanzaro et al. 2015). Therefore, to have more realistic error limits, we doubled the uncertainties of the quantities (cf. Catanzaro et al. 2015), making the $T_{\text {eff,SED }}$ uncertainties comparable to those from other photometric techniques. The results of the SED fitting procedure are listed in columns 9-12 of Table 3.

### 4.3 Atmospheric parameters from spectroscopy

In this section, we evaluate the fundamental astrophysical quantities of our targets, i.e. $T_{\text {eff }}, \log g, v \sin i$, microturbulent velocity ( $\xi$ ), and radial velocity ( $v_{\mathrm{rad}}$ ).

The first step is to determine $v_{\text {rad }}$ by computing the crosscorrelation function (CCF) using different pre-selected masks built from line lists using the code ISPEC (Blanco-Cuaresma et al. 2014; Blanco-Cuaresma 2019). We computed $v \sin i$ by comparing the observed spectra with a grid of synthetic spectra, pre-computed with plane-parallel ATLAS9 ${ }^{4}$ model atmospheres (Castelli \& Kurucz 2003) for different values of $T_{\text {eff }}$ and of $\log g$, in the wavelength range $516-519 \mathrm{~nm}$ (Mg I triplet). This comparison was performed with a least-squares method based on the minuit minimization package used by the GIRFIT package (Frémat et al. 2006).

For stars with $T_{\text {eff }} \geq 8000 \mathrm{~K}$, the Balmer lines become more sensitive to $\log g$ than to $T_{\text {eff. }}$. For such targets (KIC 4567097, KIC 4818496, KIC 5524045, KIC 5650229, KIC 6266219, KIC 7667560, and KIC 9272082), we used GIRFIT to

[^3]Table 4. List of stellar parameters from spectroscopic analysis. The different columns show: identification number, effective temperatures ( $T_{\text {eff }}$ ), surface gravity $(\log g)$, metallicity $([\mathrm{M} / \mathrm{H}])$, microturbulent velocity $(\xi)$, projected rotational velocity $(v \sin i)$, and radial velocity $\left(v_{\text {rad }}\right)$. The literature values that do not agree within $1 \sigma$ with the results from this study are given in italics.

| $\begin{aligned} & \text { KIC } \\ & \text { No. } \end{aligned}$ | $\begin{aligned} & T_{\text {eff }} \\ & (\mathrm{K}) \end{aligned}$ | $\begin{aligned} & \log g \\ & (\operatorname{dex}) \end{aligned}$ | $\begin{gathered} {[\mathrm{M} / \mathrm{H}]} \\ (\mathrm{dex}) \end{gathered}$ | $\begin{gathered} \xi \\ \left(\mathrm{km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} v \sin i \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | $\begin{gathered} v_{\mathrm{rad}} \\ \left(\mathrm{~km} \mathrm{~s}^{-1}\right) \end{gathered}$ | References |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| 3459226 | $7640 \pm 150$ | $3.94 \pm 0.14$ | $0.32 \pm 0.08$ | $3.78 \pm 0.18$ | $\begin{gathered} 73 \pm 5 \\ <120 \end{gathered}$ | $\begin{aligned} & -28.8 \pm 2.2 \\ & -24.1 \pm 24.6 \end{aligned}$ | 1 |
|  | $7410 \pm 139$ | $3.79 \pm 0.14$ |  |  |  |  | 2 |
|  | $7519 \pm 83$ | $3.62 \pm 0.16$ |  |  |  |  | 3 |
|  | $7055 \pm 140$ |  |  |  |  |  | 4 |
|  | $7516 \pm 80$ | $4.23 \pm 0.09$ |  |  |  |  | 5 |
| 4567097 | $9820 \pm 300$ | $3.90 \pm 0.12$ | $0.08 \pm 0.13$ | $1.64 \pm 0.35$ | $83 \pm 5$ | $7.1 \pm 2.4$ | 1 |
|  | $9935 \pm 360$ | $4.21 \pm 0.20$ |  |  |  |  | 3 |
|  | $8712 \pm 980$ |  |  |  |  |  | 4 |
|  | $9928 \pm 347$ |  |  |  |  |  | 5 |
| 4818496 | $9340 \pm 180$ | $3.91 \pm 0.13$ | $0.25 \pm 0.10$ | $2.26 \pm 0.23$ | $180 \pm 8$ | $11.8 \pm 5.4$ | 1 |
|  | $8941 \pm 323$ | $3.69 \pm 0.33$ |  |  |  |  | 3 |
|  | $9000 \pm 185$ |  |  |  |  |  | 4 |
|  | $8936 \pm 312$ | $3.92 \pm 0.08$ |  |  |  |  | 5 |
|  |  |  |  |  |  | $12.8 \pm 0.8$ | 6 |
| 5524045 | $9130 \pm 210$ | $3.77 \pm 0.10$ | $-0.15 \pm 0.09$ | $2.16 \pm 0.17$ | $\begin{aligned} & 218 \pm 15 \\ & 201 \pm 58 \end{aligned}$ | $3.6 \pm 13.0$ | 1 |
|  | $9267 \pm 169$ | $3.87 \pm 0.11$ |  |  |  | $7.8 \pm 36.2$ | 2 |
|  | $9217 \pm 82$ | $3.87 \pm 0.16$ |  |  |  |  | 3 |
|  | $9189 \pm 497$ |  |  |  |  |  | 4 |
|  | $9235 \pm 80$ | $3.99 \pm 0.06$ |  |  |  |  | 5 |
|  | $9500 \pm 200$ | $4.0 \pm 0.20$ |  | $1.0 \pm 0.4$ | $215 \pm 7$ |  | 7 |
| 5650229 | $8950 \pm 180$ | $3.81 \pm 0.12$ | $0.19 \pm 0.10$ | $2.02 \pm 0.19$ | $251 \pm 22$ | $-9.2 \pm 1.2$ | 1 |
|  | $9898 \pm 420$ | $3.88 \pm 0.12$ |  |  | $243 \pm 33$ | $-9.3 \pm 22.1$ | 2 |
|  | $8879 \pm 345$ | $3.87 \pm 0.24$ |  |  |  |  | 3 |
|  | $8533 \pm 400$ |  |  |  |  |  | 4 |
|  | $8842 \pm 309$ | $3.89 \pm 0.08$ |  |  |  |  | 5 |
| 6266219 | $8030 \pm 250$ | $4.11 \pm 0.17$ | $0.63 \pm 0.11$ | $3.88 \pm 0.23$ | $62 \pm 3$ | $2.0 \pm 1.5$ | 1 |
|  | $7996 \pm 282$ | $4.05 \pm 0.16$ |  |  |  |  | 3 |
|  | $7765 \pm 130$ |  |  |  |  |  | 4 |
|  | $8007 \pm 280$ | $4.09 \pm 0.09$ |  |  |  |  | 5 |
| 7667560 | $8330 \pm 200$ | $4.04 \pm 0.16$ | $0.16 \pm 0.09$ | $2.41 \pm 0.29$ | $147 \pm 14$ | $-20.2 \pm 4.8$ | 1 |
|  | $8675 \pm 320$ | $4.08 \pm 0.15$ |  |  |  |  | 3 |
|  | $8278 \pm 390$ |  |  |  |  |  | 4 |
|  | $8669 \pm 303$ | $4.22 \pm 0.09$ |  |  |  |  | 5 |
| 9272082 | $8170 \pm 190$ | $4.02 \pm 0.13$ | $0.17 \pm 0.10$ | $3.63 \pm 0.20$ | $78 \pm 3$ | $5.7 \pm 1.7$ | 1 |
|  | $7551 \pm 262$ | $3.91 \pm 0.12$ |  |  | $<120$ | $-20.8 \pm 18.3$ | 2 |
|  | $9077 \pm 334$ | $4.15 \pm 0.15$ |  |  |  |  | 3 |
|  | $8081 \pm 340$ |  |  |  |  |  | 4 |
|  | $9078 \pm 317$ | $4.06 \pm 0.08$ |  |  |  |  | 5 |
| 9349245 | $7830 \pm 260$ | $4.06 \pm 0.11$ | $0.37 \pm 0.12$ | $3.33 \pm 0.14$ | $81 \pm 3$ | $-36.2 \pm 6.3$ | 1 |
|  | $7683 \pm 79$ | $3.87 \pm 0.11$ |  |  | $<120$ | $-15.0 \pm 23.6$ | 2 |
|  | $7914 \pm 279$ | $3.72 \pm 0.29$ |  |  |  |  | 3 |
|  | $7870 \pm 393$ |  |  |  |  |  | 4 |
|  | $7913 \pm 276$ | $4.17 \pm 0.09$ |  |  |  |  | 5 |
|  | $8300 \pm 200$ | $4.00 \pm 0.30$ |  | $3.1 \pm 0.5$ | $80 \pm 3$ |  | 8 |
|  | $8000 \pm 200$ | $3.7 \pm 0.1$ |  | $2.9 \pm 0.1$ | $82 \pm 2$ |  | 9 |

Note. References: (1) This study; (2) Frasca et al. (2016); (3) Mathur et al. (2017); (4) Andrae et al. (2018); (5) Murphy et al. (2019); (6) Gontcharov (2006); (7) Niemczura et al. (2015); (8) Catanzaro et al. (2015); and (9) Niemczura et al. (2017).
derive the temperature from the spectral region with a high number of metal lines ( $495-570 \mathrm{~nm}$ ). The SNR of this spectral region was first improved using the least-squares deconvolution (LSD) method described in Tkachenko et al. (2013). Using GIRFIT, we fitted $10-15 \mathrm{~nm}$ segments of the denoised spectrum and calculated the average $T_{\text {eff }}$ for all the segments. For KIC 3459226 and KIC 9349245 ( $T_{\text {eff }}<8000 \mathrm{~K}$ ), we determined $T_{\text {eff }}$ using hydrogen line profiles. The errors are a con-
tribution of uncertainties encountered in continuum normalization and errors from the fitting procedure. We adopted an error of 100 K as a contribution from continuum normalization (Catanzaro et al. 2015).

The values of $\log g$ were estimated by two different methods depending on $T_{\text {eff }}$ of each object. For target stars with $T_{\text {eff }}>$ 8000 K , we computed $\log g$ from fitting hydrogen line wings. At $T_{\text {eff }}<8000 \mathrm{~K}$, the Balmer lines lose their sensitivity to


Figure 3. The $\mathrm{H} \beta$ (left), Mg I triplet (middle), and $\mathrm{H} \alpha$ (right) spectral regions for the target stars. The observed (black) and synthetic spectra (red) are shown. The root mean square value of the residuals (Res.rms) for each spectrum is given in the bottom right corner of each $\mathrm{H} \beta$ panel.


Figure 4. The $T_{\text {eff }}$ and $\log g$ from the current study compared with those reported by Mathur et al. (2017) (red), Andrae et al. (2018) (black), and Murphy et al. (2019) (blue). The respective correlation coefficient, $r$, with the weighted average of the literature values is given in the top right corner of each panel.
$\log g$, and so metal lines (Fe I/Fe II and/or $\mathrm{MgI}_{\text {I }}$ triplet) were used instead.

We uniquely obtained $\xi$ by fitting iron lines using the MOOG radiative transfer code (Sneden et al. 2012), ATLAS9 model atmospheres, the Vienna Atomic Line Database (VALD) line list (Kupka et al. 1999), and Asplund et al. (2009) solar abundances, all incorporated in an integrated software ISPEC. We determined the initial guess for $\xi$ from the relation (Gebran et al. 2014)
$\xi=3.31 \times \exp \left[-\left(\log \left(\frac{T_{\mathrm{eff}}}{8071.03}\right)^{2} / 0.01045\right)\right]$.
Fig. 3 shows the $\mathrm{H} \beta$ (left-hand panels), Mg I triplet (middle panels), and $\mathrm{H} \alpha$ (right-hand panels) spectral regions for the target stars (black) and the synthetic spectra (red) computed with the final results of the atmospheric parameters. The results for the atmospheric fundamental quantities and their errors are listed in Table 4. The spectroscopic $T_{\text {eff }}$ values differ from those estimated using photometric indices ( $u b v y \beta$, 2MASS, and UBV) and SED by an


Figure 5. The projected rotational velocity $v \sin i$ as computed in this study compared with the equatorial rotational velocity $v_{\text {rot }}$ by Trust et al. (2020). The blue, red, and magenta dots are normal A, Am, and marginal Am stars, respectively. The blue dashed line represents $\sin i=1$, while the green and black ones correspond to $i=60^{\circ}$ and $30^{\circ}$, respectively.
average temperature of about $360,180,350$, and 180 K , respectively. The spectroscopic $\log g$ and $[\mathrm{M} / \mathrm{H}]$ values are more consistent with those estimated from ubvy $\beta$ photometry (average $\Delta \log g_{u v b y \beta}=$ 0.18 dex and $\Delta[\mathrm{M} / \mathrm{H}]_{\text {uvby } \beta}=0.10$ dex) than those from SED (average $\Delta \log g_{\text {SED }}=0.5$ dex and $\Delta[\mathrm{M} / \mathrm{H}]_{\text {SED }}=0.23$ dex $)$. Our final results (from spectroscopy) agree in general within the error limits with those from previous studies. The literature values for which this is not the case are shown in italics in Table 4. Particularly, we compared our $T_{\text {eff }}$ and $\log g$ with those computed by Mathur et al. (2017), Andrae et al. (2018), and Murphy et al. (2019) as shown, in panels (a) and (b) of Fig. 4, respectively. Based on the correlation coefficient (presented in the top right corner of each panel) with the weighted average of the literature values, $T_{\text {eff }}$ is more consistent than $\log g$. The $v \sin i$ values are excellently consistent with the equatorial rotational velocity $\left(v_{\text {rot }}\right)$ computed by Trust et al. (2020) as shown in Fig. 5, where all the stars are positioned either on or below the $\sin i=1$ line.

### 4.4 Individual chemical abundances

The individual chemical abundances were determined based on direct fitting of the theoretical profiles of individual spectral lines. Since some targets are fast rotators and others are CP stars, lines are often blended and direct spectrum synthesis takes into account possible blends (Romanovskaya et al. 2019). We used the SYNTHV_NLTE code (Tsymbal, Ryabchikova \& Sitnova 2019) and a grid of atmosphere models pre-computed with the LLMODELS code (Shulyak et al. 2004). This was implemented through an IDL visualization program, BINMAG6 (Kochukhov 2018). During spectral synthesis, local thermodynamic equillibrium (LTE) was assumed. For oxygen abundance analysis, we excluded the O I triplet (around 777.1-777.5 nm), since it is affected by non-LTE effects (Przybilla et al. 2000). For the same reason, we never considered Na I lines at 589.0 and 589.6 nm (Takeda et al. 2009) in our analysis. We adopted spectral line lists and atomic
Table 5. Individual chemical abundances inferred for our target stars. The solar abundances are listed in the last column (Asplund et al. 2009). The values are expressed in the form $\log \left(N_{\mathrm{el}} / N_{\text {Tot }}\right)$. The number of segments from which each element was determined is given in brackets.

|  | KIC 3459226 | KIC 4567097 | KIC 4818496 | KIC 5524045 | KIC 5650229 | KIC 6266219 | KIC 7667560 | KIC 9272082 | KIC 9349245 | Solar |
| :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: | :---: |
| C | $-3.93 \pm 0.16$ (3) | $-4.08 \pm 0.11$ (3) | $-3.63 \pm 0.10$ (4) | $-3.77 \pm 0.09$ (6) | $-3.87 \pm 0.10$ (6) | $-3.63 \pm 0.12$ (9) | $-4.16 \pm 0.09$ (5) | $-3.87 \pm 0.13$ (5) | $-3.75 \pm 0.09$ (6) | -3.57 |
| O | $-3.10 \pm 0.08^{a}(1)$ | $-3.22 \pm 0.15^{a}(1)$ | $-2.48 \pm 0.09^{a}(1)$ | $-2.61 \pm 0.10^{a}$ (2) | $-3.22 \pm 0.16^{a}$ (2) | $-3.06 \pm 0.10^{a}(2)$ | $-3.01 \pm 0.15^{a}(1)$ | $-2.95 \pm 0.09^{a}(1)$ | $-2.60 \pm 0.10^{a}(1)$ | -3.31 |
| Na | $-4.12 \pm 0.07^{a}(1)$ | $-4.01 \pm 0.20^{a}(1)$ | $-5.02 \pm 0.10^{a}(1)$ | - | - | $-4.75 \pm 0.11^{a}$ (2) | $-4.92 \pm 0.13^{a}(2)$ | $-5.01 \pm 0.15^{a}(2)$ | $-4.89 \pm 0.09^{a}(2)$ | -5.76 |
| Mg | $-4.23 \pm 0.20$ (6) | $-4.30 \pm 0.14$ (6) | $-3.95 \pm 0.11$ (7) | $-4.47 \pm 0.10$ (8) | $-4.19 \pm 0.22$ (4) | $-4.08 \pm 0.11$ (9) | $-4.31 \pm 0.13$ (5) | $-4.35 \pm 0.11$ (10) | $-4.08 \pm 0.09$ (9) | -4.40 |
| Al | $-4.61 \pm 0.09^{a}(1)$ | - | - | - | $-4.98 \pm 0.11^{a}(1)$ | - | - | $-5.25 \pm 0.13^{a}(1)$ | - | -5.55 |
| Si | $-3.91 \pm 0.16$ (8) | $-4.27 \pm 0.14^{a}(2)$ | $-4.00 \pm 0.11^{a}(1)$ | $-4.01 \pm 0.11$ (3) | $-3.71 \pm 0.15$ (3) | $-3.82 \pm 0.11$ (10) | $-3.83 \pm 0.17^{a}(2)$ | $-4.40 \pm 0.09$ (12) | $-3.73 \pm 0.10$ (5) | - 4.49 |
| S | $-4.22 \pm 0.10^{a}(1)$ | - | - | - | - | $-4.29 \pm 0.09$ (3) | - | $-4.48 \pm 0.15^{\text {a }}$ (1) |  | -4.88 |
| K | $-5.92 \pm 0.09^{a}(1)$ | - | - | - | - | $-6.63 \pm 0.09^{a}(1)$ | - | $-6.97 \pm 0.09^{a}(1)$ | - | -6.97 |
| Ca | $-6.05 \pm 0.14$ (11) | $-5.89 \pm 0.12$ (3) | $-5.18 \pm 0.14$ (3) | $-5.78 \pm 0.18$ (5) | $-5.79 \pm 0.12$ (5) | $-5.67 \pm 0.10$ (23) | $-5.48 \pm 0.20$ (6) | $-5.68 \pm 0.09$ (18) | $-5.38 \pm 0.10$ (20) | -5.66 |
| Sc | $-9.35 \pm 0.12$ (4) | $-8.92 \pm 0.17$ (5) | $-9.04 \pm 0.12$ (4) | $-9.06 \pm 0.10$ (4) | $-8.82 \pm 0.24$ (4) | $-9.11 \pm 0.16$ (7) | $-8.66 \pm 0.13$ (5) | $-8.70 \pm 0.10$ (10) | $-8.96 \pm 0.09$ (3) | -8.85 |
| Ti | $-6.88 \pm 0.17$ (30) | $-7.36 \pm 0.10$ (12) | $-6.87 \pm 0.12$ (11) | $-7.22 \pm 0.10$ (19) | $-7.16 \pm 0.16$ (15) | $-6.57 \pm 0.09$ (24) | $-6.96 \pm 0.13$ (26) | $-6.97 \pm 0.12$ (27) | $-6.77 \pm 0.10$ (23) | -7.05 |
| V | $-7.80 \pm 0.10^{a}(2)$ | $-7.87 \pm 0.13^{a}(1)$ | $-7.85 \pm 0.15^{a}(1)$ | - | $-7.86 \pm 0.12^{a}(2)$ | $-6.98 \pm 0.12$ (4) | $-8.22 \pm 0.14^{a}(2)$ | $-8.01 \pm 0.10$ (3) | $-7.32 \pm 0.12$ (5) | -8.07 |
| Cr | $-6.04 \pm 0.13$ (5) | $-6.58 \pm 0.09$ (7) | $-6.39 \pm 0.09$ (7) | $-6.95 \pm 0.23$ (7) | $-6.30 \pm 0.10$ (7) | $-5.53 \pm 0.11$ (17) | $-6.36 \pm 0.13$ (10) | $-6.30 \pm 0.09$ (19) | $-6.15 \pm 0.09$ (18) | -6.36 |
| Mn | $-6.16 \pm 0.19$ (3) | $-6.37 \pm 0.11^{a}(1)$ | $-6.38 \pm 0.12^{a}(1)$ | $-6.74 \pm 0.12$ (3) | - | $-6.31 \pm 0.10$ (3) | $-6.89 \pm 0.12$ (3) | $-7.02 \pm 0.23$ (5) | $-6.18 \pm 0.08$ (5) | -6.57 |
| Fe | $-3.89 \pm 0.11$ (79) | $-4.63 \pm 0.11$ (17) | $-4.24 \pm 0.13$ (38) | $-4.68 \pm 0.11$ (33) | $-4.27 \pm 0.13$ (31) | $-3.82 \pm 0.10$ (83) | $-4.25 \pm 0.14$ (43) | $-4.61 \pm 0.09$ (43) | $-4.18 \pm 0.09$ (44) | -4.50 |
| Co | $-5.94 \pm 0.09$ (6) | - | - | - | $-6.34 \pm 0.11^{a}(1)$ | $-5.71 \pm 0.13$ (8) | - | $-6.67 \pm 0.16$ (3) | - | -7.01 |
| Ni | $-5.10 \pm 0.10$ (16) | $-5.46 \pm 0.14^{a}(1)$ | $-6.22 \pm 0.19^{a}(1)$ | $-5.23 \pm 0.09^{a}(1)$ | $-4.67 \pm 0.16$ (21) | $-5.14 \pm 0.09$ (21) | $-5.37 \pm 0.15$ (6) | $-5.88 \pm 0.09$ (9) | $-5.29 \pm 0.09$ (9) | - 5.78 |
| Cu | $-7.05 \pm 0.13$ (3) | - | - | - | - | $-6.37 \pm 0.23$ (3) | - | $-7.17 \pm 0.21$ (3) | - | - 7.81 |
| Zn | $-7.02 \pm 0.09$ (3) | - | - | - | - | $-7.27 \pm 0.13^{a}(1)$ | - | $-7.29 \pm 0.09^{a}(1)$ | $-7.17 \pm 0.09^{a}(1)$ | -7.44 |
| Sr | $-8.12 \pm 0.11$ (4) | $-9.36 \pm 0.12^{a}(2)$ | - | $-9.04 \pm 0.15^{a}(1)$ | $-8.84 \pm 0.10^{a}(1)$ | $-8.29 \pm 0.09^{a}(2)$ | $-8.84 \pm 0.22^{a}$ (2) | $-8.75 \pm 0.15^{a}(1)$ | $-8.49 \pm 0.14$ (3) | -9.13 |
| Y | $-8.83 \pm 0.19$ (6) | - | $-9.81 \pm 0.11^{a}(1)$ | $-9.83 \pm 0.11^{a}(1)$ | - | $-9.03 \pm 0.19$ (6) | $-9.78 \pm 0.09$ (3) | $-9.62 \pm 0.11$ (6) | $-9.42 \pm 0.09$ (5) | - 9.79 |
| Zr | $-8.51 \pm 0.21$ (4) | - | - | - | $-9.51 \pm 0.21$ (3) | $-8.98 \pm 0.10$ (3) | $-9.09 \pm 0.11^{a}(1)$ | $-9.18 \pm 0.14^{a}(2)$ | $-9.12 \pm 0.17$ (13) | - 9.47 |
| Ba | $-8.69 \pm 0.19$ (5) | - | - | - | $-10.06 \pm 0.11^{a}$ (2) | $-8.44 \pm 0.17$ (5) | $-9.77 \pm 0.16$ (3) | $-9.45 \pm 0.14$ (5) | $-8.76 \pm 0.14$ (3) | -9.82 |
| La | $-8.61 \pm 0.13$ (8) | - | - | - | - | $-8.91 \pm 0.21$ (10) | - | $-10.63 \pm 0.10$ (4) | - | - 10.90 |
| Ce | $-8.93 \pm 0.18$ (7) | - | - | - | - | $-8.61 \pm 0.11$ (9) | - | $-10.02 \pm 0.12$ (5) | - | - 10.42 |
| Pr | $-9.79 \pm 0.10^{a}$ (1) | - | - | - | - | $-9.07 \pm 0.10^{a}(2)$ | - | - | - | - 11.28 |
| Nd | $-8.73 \pm 0.11$ (5) | - | - | - | - | $-8.79 \pm 0.18$ (7) | - | - | - | - 10.58 |

[^4]

Figure 6. The individual chemical abundance patterns for our targets. The horizontal dashed line represents solar abundance (Asplund et al. 2009). The grey background highlights the light elements ( Ca and Sc ) and heavy elements ( $\mathrm{Sr}, \mathrm{Y}, \mathrm{Zr}, \mathrm{Ba}, \mathrm{La}, \mathrm{Ce}, \mathrm{Pr}$, and Nd ) important for Am classification.


Figure 7. A comparison of our individual chemical abundances for KIC 5524045 (top panel) with values by Niemczura et al. (2015) and KIC 9349245 (bottom panel) with those determined by Catanzaro et al. (2015). The corresponding correlation coefficient, $r$, is given in the top right corner of each panel.
parameters from the 3D release of the Vienna Atomic Line Database (VALD3; Ryabchikova et al. 2015).

We divided each spectrum into several intervals, each $\approx 5 \mathrm{~nm}$ wide, and derived the abundances in each interval by performing an $\chi^{2}$ minimization of the difference between the observed and synthetic spectrum. The average individual chemical abundances are listed in Table 5. We used the errors on $T_{\text {eff }}, \log g, v \sin i, \xi$, position of the continuum, and accuracy of the oscillator strengths $(\log (\mathrm{gf}))$ of the considered lines to derive abundance uncertainties. For each element, the errors on abundance associated with the errors on each of the $T_{\text {eff }}$, $\log g, v \sin i, \xi$, position of the continuum, and the $\log$ (gf) were added together in quadrature to get the total uncertainty. We note that the abundances of elements whose lines only appear in less than three segments should be taken with caution since they were computed from few lines.

Table 6. The different columns show: tag (for identification in the HR diagram in Fig. 9), identification (KIC) number, absolute magnitude $\left(M_{\mathrm{V}}\right)$, bolometric correction $(\mathrm{BC})$, luminosity $\left(\log \left(L_{\star} / L_{\odot}\right)\right.$ ), and stellar radius $(R)$ as estimated from standard relations.

| Tag | KIC <br> No. | $M_{\mathrm{v}}$ <br> $(\mathrm{mag})$ | BC <br> $(\mathrm{mag})$ | $\log \left(L_{\star} / L_{\odot}\right)$ <br> $(\mathrm{dex})$ | $R$ <br> $\left(R_{\odot}\right)$ |
| :--- | :---: | ---: | :---: | :---: | :---: |
| a | 3459226 | $2.616 \pm 0.145$ | $0.0304 \pm 0.0038$ | $0.846 \pm 0.056$ | $1.50 \pm 0.15$ |
| b | 4567097 | $-0.820 \pm 0.132$ | $-0.2149 \pm 0.0610$ | $2.319 \pm 0.028$ | $4.98 \pm 0.46$ |
| c | 4818496 | $0.935 \pm 0.092$ | $-0.1735 \pm 0.0410$ | $1.600 \pm 0.020$ | $2.28 \pm 0.15$ |
| d | 5524045 | $0.918 \pm 0.116$ | $-0.0743 \pm 0.0303$ | $1.567 \pm 0.034$ | $2.46 \pm 0.20$ |
| e | 5650229 | $0.596 \pm 0.128$ | $-0.0512 \pm 0.0276$ | $1.687 \pm 0.040$ | $2.92 \pm 0.26$ |
| f | 6266219 | $1.848 \pm 0.122$ | $0.0239 \pm 0.0024$ | $1.156 \pm 0.050$ | $1.98 \pm 0.24$ |
| g | 7667560 | $1.850 \pm 0.108$ | $0.0140 \pm 0.0130$ | $1.159 \pm 0.038$ | $1.83 \pm 0.17$ |
| h | 9272082 | $0.924 \pm 0.105$ | $0.0225 \pm 0.0047$ | $1.526 \pm 0.040$ | $2.93 \pm 0.27$ |
| i | 9349245 | $2.290 \pm 0.083$ | $0.0279 \pm 0.0084$ | $0.978 \pm 0.030$ | $1.71 \pm 0.17$ |




Figure 8. Stellar radius and $\log \left(L_{\star} / L_{\odot}\right)$ from the current study compared with those determined by Murphy et al. (2019). The respective correlation coefficient, $r$, is given in the top right corner of each panel.


Figure 9. The HR diagram showing evolutionary phase of target stars. The blue, red, and magenta dots are normal A, Am, and marginal Am stars, respectively. The black solid line represents the zero-age main sequence (ZAMS) while the green solid lines are evolutionary tracks for different masses. The blue and red solid lines indicate the blue and red edges of the observational $\delta$ Scuti instability strip reported by Murphy et al. (2019).

## 5 COMMENTS ON INDIVIDUAL STARS

In this section, we present the findings of the abundance analysis obtained for each target. The derived abundances and the estimated uncertainties are expressed as $\log \left(N_{\mathrm{el}} / N_{\text {Tot }}\right)$ and are reported in Table 5. Fig. 6 shows the abundance patterns for each star relative to the solar abundances (Asplund et al. 2009).

### 5.1 KIC 3459226 (BD + 38 3705)

This star was classified as kA2hF0mF2 by Frasca et al. (2016) using LAMOST data (Cui et al. 2012; Zhao et al. 2012). We classified this target as kA 2 hFO mF 3 using MKCLASS with HERMES data. Calcium and scandium are $\approx 0.4$ and $\approx 0.5$ dex below solar values, respectively, while the iron-peak and heavy elements such as $\mathrm{Sr}, \mathrm{Y}$, $\mathrm{Zr}, \mathrm{Ba}, \mathrm{La}, \mathrm{Ce}, \mathrm{Pr}$, and Nd are clearly overabundant (values up to
$\approx 2$ dex). These overabundances and spectral classification confirm the Am peculiarity of this star.

### 5.2 KIC 4567097 (HD 184469)

KIC 4567097 was reported in the Henry Draper (HD) catalogue (Cannon \& Pickering 1993) as a B9 star. Using mKCLASS, we obtained a spectral type of B9 III. No abundance analysis was found in the literature. The calcium abundance is $\approx 0.23$ dex below solar value while scandium and the iron-peak elements are generally solarlike. Only few absorption lines were available and used to determine abundances. Based on the spectral class and chemical abundances obtained, we propose to consider KIC 4567097 as a non-Am star.

### 5.3 KIC 4818496 (HD 177592)

This object is listed in the HD catalogue as an A0 star (Cannon \& Pickering 1993) while our analysis classifies it as A1 V. The $T_{\text {eff }}$ in this study is slightly higher than the values in the literature while our $\log g$ value is consistent (see Table 4). We found this star to be a fast rotator with a $v \sin i$ value of $180 \pm 8 \mathrm{~km} \mathrm{~s}^{-1}$, which is consistent with the value computed by Trust et al. (2020). O, Na, $\mathrm{Mg}, \mathrm{Si}$, and Ca are overabundant ( $0.5-0.9 \mathrm{dex}$ ) while Sc and iron-peak elements are, on average, of solar content. Based on these results, we classify this star as a (chemically) normal A-type star.

### 5.4 KIC 5524045 (BD +40 3639)

KIC 5524045 was first classified as A0.5 Va+ ('a+' denotes higher luminosity MS star) by Niemczura et al. (2015) and subsequently as A1 V by Frasca et al. (2016). They are both in excellent agreement with our classification as a 1 V star. We observed this star to be fast rotating with $v \sin i=219 \pm 12 \mathrm{~km} \mathrm{~s}^{-1}$. Similar rotational velocities, $215 \pm 7 \mathrm{~km} \mathrm{~s}^{-1}$ (Niemczura et al. 2015) and $201 \pm 58 \mathrm{~km} \mathrm{~s}^{-1}$ (Frasca et al. 2016) were published before. All the measured elements are on average of solar type. The top panel of Fig. 7 shows that our abundances are in agreement with those determined by Niemczura et al. (2015). Based on the abundance pattern, we classify this star as normal.

### 5.5 KIC 5650229 (HD 226697)

This object was first listed as an A2 star in the HD catalogue (Cannon \& Pickering 1993). Recently, Frasca et al. (2016) classified it as a B9 III or IV star. Our analysis indicates a spectral type as A1 IV. The most recent study (Murphy et al. 2019) presents the star's $T_{\text {eff }}$ and $\log g$ as $8842 \pm 309 \mathrm{~K}$ and $3.89 \pm 0.08 \mathrm{~cm} \mathrm{~s}^{-2}$, respectively, which agrees with our result $\left(8920 \pm 190 \mathrm{~K}\right.$ and $\left.3.82 \pm 0.11 \mathrm{~cm} \mathrm{~s}^{-2}\right)$. This star is a fast rotator whose $v \sin i$ was first reported as $243 \pm 33 \mathrm{~km} \mathrm{~s}^{-1}$ (Frasca et al. 2016), which is consistent with our outcome of $252 \pm 24 \mathrm{~km} \mathrm{~s}^{-1}$. With exception of $\mathrm{Al}, \mathrm{Si}, \mathrm{Co}$, and Ni that are overabundant (with an average abundance of $\approx 0.8 \mathrm{dex}$ ), the abundances of the other elements are solar-like and hence we consider this star as non-Am.

### 5.6 KIC 6266219 (TYC 3127-2016-1)

KIC 6266219 was listed among A-type stars in the nominal Kepler field (Balona 2013). We obtained its spectral type as kA3hA7mF1. The available $T_{\text {eff }}=7996 \pm 282 \mathrm{~K}$ and $\log g=4.05 \pm 0.16 \mathrm{~cm} \mathrm{~s}^{-2}$ (Mathur et al. 2017), $T_{\text {eff }}=7765 \pm 130 \mathrm{~K}$ (Andrae et al. 2018), and $T_{\text {eff }}=8007 \pm 280 \mathrm{~K}$ and $\log g=4.09 \pm 0.09 \mathrm{~cm} \mathrm{~s}^{-2}$ (Murphy
et al. 2019) are consistent with our results. This object is a moderate rotator with $v \sin i$ of $60 \mathrm{~km} \mathrm{~s}^{-1}$, which agrees with the value given by Trust et al. (2020). The chemical abundances of Ca and Sc are solar-like while iron-peak and heavy elements are overabundant up to $\approx 2$ dex. Based on the spectral type and chemical abundance pattern, we classify KIC 6266219 as an Am star.

### 5.7 KIC 7667560 (HD 177458)

KIC 7667560 was listed in the HD catalogue as an A-type star by Cannon \& Pickering (1993). Our spectral classification analysis revealed a spectral type of A3 IV. The $T_{\text {eff }}$ and $\log g$ values published by Mathur et al. (2017), Andrae et al. (2018), and Murphy et al. (2019) and are fully consistent with the values obtained in this study. $\mathrm{Ca}, \mathrm{Sc}$, iron-peak, and heavy elements, on average, have solar-like abundances. $\mathrm{Na}(\approx 0.83 \mathrm{dex})$ and $\mathrm{Si}(\approx 0.66 \mathrm{dex})$ are overabundant. The chemical abundance pattern does not show Am characteristics.

### 5.8 KIC 9272082 (HD 179458)

This star was first classified as Am by Macrae (1952) and subsequently by Bertaud (1960). The most recent classification as A3 V, by Frasca et al. (2016), matches our result, A4 V. As shown in Table 4, only the $T_{\text {eff }}$ computed by Andrae et al. (2018) using Gaia observations agrees with the $T_{\text {eff }}$ obtained in this study. Our $\log g$ value agrees well with the values: $3.91 \pm 0.12$ dex (Frasca et al. 2016), $4.15 \pm 0.15$ dex (Mathur et al. 2017), and $4.06 \pm 0.08$ dex (Murphy et al. 2019). With exception of $\mathrm{Na}(\approx 0.75$ dex), which is overabundant, abundances of the rest of the elements are solar-like. Based on the spectral classification and the chemical abundance pattern, we consider this star as non-Am.

### 5.9 KIC 9349245 (HD 185658)

KIC 9349245 is an MS A-type star as reported in various studies (e.g. Cannon \& Pickering 1993; Frasca et al. 2016; Gray et al. 2016) and it is included in the catalogue of CP stars (Renson \& Manfroid 2009). We determined its spectral type as A8 IV (Sr), which is similar to the one given by Gray et al. (2016). The atmospheric parameters and velocities ( $\xi, v \sin i$ and $v_{\text {rad }}$ ), within error limits, generally agree with values in the literature except for $T_{\text {eff }}$ of Catanzaro et al. (2015) (see Table 4). The Ca and Sc abundances are solar-like, while ironpeak and heavy metals are moderately overabundant (with an average abundance of $\approx 0.5$ dex). The chemical abundance pattern implies a marginal Am star. Similar observations were made by Catanzaro et al. (2015). The bottom panel of Fig. 7 shows a good agreement between the two studies.

## 6 LOCATION IN THE HR DIAGRAM

A precise positioning of both Am and non-Am stars in the HR diagram is useful for investigating possible systematic differences in the region occupied by these objects with regard to normal A stars. Moreover, all the targets studied in this paper are 'hump and spike' stars and their position is useful to constrain the location of stars with $r$ modes in the HR diagram.

The luminosity $\left(\log \left(L_{\star} / L_{\odot}\right)\right)$ of the target stars was calculated using the standard technique. We computed extinction from $A_{\mathrm{V}}=$ $3.1 \times E(\mathrm{~B}-\mathrm{V})$, where $E(B-V)$ was obtained as discussed in Section 4.1. The bolometric correction (BC) was calculated using the temperature-dependent function of Flower (1996) after a revision by

Torres (2010), while the absolute magnitude ( $M_{\mathrm{v}}$ ) was determined using the Gaia parallaxes (Gaia Collaboration 2018). The uncertainties in $\log \left(L_{\star} / L_{\odot}\right)$ result from uncertainties in $V$-band indices, parallax, and $T_{\text {eff. }}$. Having obtained precise $T_{\text {eff }}$ and $\log \left(L_{\star} / L_{\odot}\right)$ values, we also determined the stellar radius using the standard relation (from the Stefan-Boltzmann law) (Boltzmann 1884; Paul et al. 2015; Montambaux 2018),
$\log \left(\frac{R}{R_{\odot}}\right)=0.5 \log \left(\frac{L_{\star}}{L_{\odot}}\right)-2.0 \log \left(T_{\text {eff }}\right)+7.52340473$.
The results and their uncertainties are listed in Table 6. Our values for $R$ (except for KIC 4818496 and KIC 9272082) and $\log \left(L_{\star} / L_{\odot}\right)$ are in good agreement with those determined by Murphy et al. (2019), as shown in panels (a) and (b) of Fig. 8, respectively.

Fig. 9 is the HR diagram showing the location of the normal A (blue) and Am (red) stars. With exception of KIC 4567097, which could be on the red giant phase, all our targets are located in the MS. This was also observed for KIC 4567097 from MK spectral classification in Section 3, where the luminosity class is III (giant), and from the spectral type obtained from SED fitting as discussed in Section 4.2. KIC 9272082 could be close to the end of its MS lifetime. All the Am stars lie within the observational $\delta$ Scuti instability strip determined by Murphy et al. (2019). With exception of KIC 7667560 and KIC 9272082, the normal A-type stars lie outside this instability strip.

## 7 CONCLUSION

In this paper, we present a spectroscopic analysis of a sample of nine objects that are classified in literature as 'hump and spike' stars. The sample comprises of both normal A and Am stars. The analysis is based on high-resolution spectra obtained with HERMES mounted at the Cassegrain focus of the $1.2-\mathrm{m}$ Mercator telescope located at La Palma, Spain. For each star, we determined the spectral type and obtained fundamental parameters such as the effective temperature $T_{\text {eff }}$, surface gravity $\log g$, metallicity [M/H], microturbulent velocity $\xi$, projected rotational velocity $v \sin i$, and radial velocity $v_{\text {rad }}$. We also performed a detailed individual chemical abundance analysis for each target, obtained their luminosity $\log \left(L_{\star} / L_{\odot}\right)$ and derived an accurate position in the HR diagram.

All the stars are either fast (KIC 4818496, KIC 5524045, KIC 5650229, and KIC 7667560) or moderate (KIC 3459226, KIC 4567097 , KIC 6266219 , KIC 9272082 , and KIC 9349245 ) rotators leading to severe blending of their spectral lines. To overcome this problem, we applied the spectral synthesis method in our analysis using different codes. Of all the stars in our sample, only two stars have high-resolution spectroscopic analysis reported in the literature, i.e. KIC 5524045 (Niemczura et al. 2015) and KIC 9349245 (Catanzaro et al. 2015). For the remaining seven targets, it was the first time that a detailed spectroscopic analysis was done.

We classified KIC 4567097 as a giant while the rest are dwarfs. So far, KIC 6266219 and KIC 9272082 have been considered as chemically normal and Am, respectively (e.g. Trust et al. 2020). Based on our spectral classification analysis and the resulting chemical abundance pattern, we reclassify KIC 6266219 as an Am star (kA3hA7mF1) while KIC 9272082 as chemically normal. Therefore, we conclude that our sample contains two Am stars (KIC 3459226 and KIC 6266219) and one marginal Am star (KIC 9349245), and six targets that are considered as non-Am stars (KIC 4567097, KIC 4818496, KIC 5524045, KIC 5650229, KIC 7667560, and KIC 9272082). With exception of the moderate rotators KIC 4567097 and KIC 9272082 , the remaining four non-

Am stars are fast rotators ( $v \sin i>120 \mathrm{~km} \mathrm{~s}^{-1}$ ), a property that facilitates mixing against atomic diffusion while minimizing the chemical peculiarity. In addition, Fig. 5 nicely shows that the Am and marginal Am stars are the slow rotators in the sample, which supports the idea that slow rotation facilitates the creation of chemical peculiarities.

The availability of HERMES and other spectrographs capable of high-resolution spectroscopy will allow us, in future, to continue the study of the 'hump and spike' stars. Particularly, we will segregate Am from non-Am stars in the 'hump and spike' star sample. In the near future, modelling of these stars will shed light on the interaction of chemical peculiarity, rotation, and Rossby waves.

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## DATA AVAILABILITY

The data underlying this article are available in the Figshare repository at https://dx.doi.org/10.6084/m9.figshare.13604348.v2.

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[^1]:    ${ }^{1}$ http://iraf.noao.edu/

[^2]:    ${ }^{2}$ https://gcpd.physics.muni.cz
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[^3]:    ${ }^{4}$ http://www.stsci.edu/hst/observatory/crds/castelli_kurucz_atlas.html

[^4]:    ${ }^{a}$ These elements were only found in less than three segments.

